

Three-dimensional Finite-element Forward Modeling of Electromagnetic Data using Potentials and Unstructured Meshes

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Introduction and Motivation

Forward Modeling

Definition:
 What a sensor (A loop of current for example) would measure if the actual distribution of the subsurface materials is given model.

Purpose:

In order to extract as much information as possible from survey data forward modeling for complex 3D Earth models are required.

Structured Vs. Unstructured Grids

Unstructured grids provide more flexibility in generating the complex-shaped Earth geometries such as curved objects (see Figure 2-b) and topography features. Furthermore from the computational standpoint, the ability to locally refine the mesh at the regions of interest and coarsening it near the transition boundaries (see Figure 2-d) is a particular benefit of using unstructured meshes. In fact, this could be an advantage over the structured rectangular grids (see Figure 2-c) where having small cells close to the boundaries of the mesh is inevitable and results in computational inefficiencies.

Inductive and Galvanic parts

Despite the long history of EM methods in geophysical exploration, there has not yet been a thorough investigation of the inductive and galvanic effects and interplay between them. Induction is characterized by coupling via a time varying magnetic field, both between a closed-loop transmitter and an isolated conductor. By contrast, a synopsis of the galvanic process would be the direct flow of the current in and from the grounded source similar to what is observed in DC resistivity methods. In fact, channeling of the current at the regions of conductivity gradient characterizes the galvanic response.

Methodology

E-Field Equation:

In the quasi-static regime the electric field satisfies the following equation:
 $\nabla \times \mathbf{E} = -i\omega \mu_0 \mathbf{J}^*$
 where \mathbf{E} is the electric field, μ_0 is the magnetic permeability of free space, σ is the electrical conductivity, and \mathbf{J}^* is the current density of the source.

Decomposition:

In one formulation the electric field is decomposed into vector magnetic potential, \mathbf{A} , (induced by \mathbf{J}) and scalar electric potential, ϕ .

Equations to be Discretized:

$$\nabla \times \nabla \times \mathbf{E} + i\omega \mu_0 \sigma \mathbf{E} = -i\omega \mu_0 \mathbf{J}^*$$

$$\mathbf{E} = -i\omega \mathbf{A} - \nabla \phi$$

$$\nabla \times \nabla \times \mathbf{A} + i\omega \mu_0 \sigma \mathbf{A} + \mu_0 \sigma \nabla \phi = \mu_0 \mathbf{J}^*$$

$$-i\omega \nabla \cdot (\sigma \mathbf{A} - \nabla \phi) = -\nabla \cdot \mathbf{J}^*$$

$$\mathbf{A} = \sum_{j=1}^{N_{nodes}} \hat{\mathbf{A}}_j N_j \quad \phi = \sum_{k=1}^{N_{nodes}} \hat{\phi}_k N_k$$

The System of Equations Solved:

$$\begin{pmatrix} \mathbf{C} & -\omega \mu_0 \mathbf{D} & \mu_0 \mathbf{F} & \mathbf{0} \\ \omega \mu_0 \mathbf{D} & \mathbf{C} & \mathbf{0} & \mu_0 \mathbf{F} \\ \mathbf{0} & -\omega \mu_0 \mathbf{F}^T & \mathbf{H} & \mathbf{0} \\ \omega \mu_0 \mathbf{F}^T & \mathbf{0} & \mathbf{0} & \mathbf{H} \end{pmatrix} \begin{pmatrix} \hat{\mathbf{A}} \\ \hat{\phi} \\ \hat{\phi}^* \\ \hat{\mathbf{A}}^* \end{pmatrix} = \begin{pmatrix} \mu_0 \mathbf{S}_1 \\ \mathbf{0} \\ \mathbf{0} \\ \mathbf{S}_2 \end{pmatrix}$$

where $\mathbf{C}_ij = \int_{\Omega} (\nabla_i \times \mathbf{N}_j) \cdot (\nabla_i \times \mathbf{N}_i) d\Omega$, $\mathbf{F}_ij = \int_{\Omega} \nabla_i N_j \cdot (\sigma \mathbf{N}_i) d\Omega$
 $\mathbf{D}_ij = \int_{\Omega} \sigma \mathbf{N}_i \cdot \mathbf{N}_j d\Omega$, \mathbf{F}^T is the transpose of the matrix \mathbf{F}
 $\mathbf{H}_{ij} = \int_{\Omega} \nabla_i N_j \cdot (\sigma \nabla_i \phi) d\Omega$ with $i, j = 1, \dots, N_{nodes}$ and $L, K = 1, \dots, N_{nodes}$
 $\hat{\mathbf{A}}^i$ and $\hat{\phi}^i$ are real and imaginary parts of the approximated vector potential \mathbf{A}^i and scalar potential ϕ^i , and $\hat{\phi}^i$ and $\hat{\mathbf{A}}^i$ are real and imaginary parts of the approximated scalar potential ϕ^i and vector potential \mathbf{A}^i respectively.

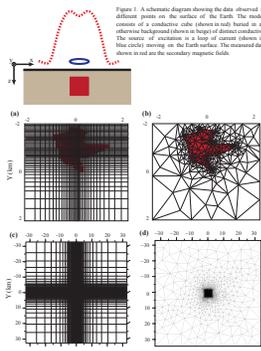


Figure 2. Generating the current independent body using (a) structured rectangular and (b) unstructured tetrahedral grids. An inset view of the mesh subdivided into (c) structured and (d) unstructured elements.

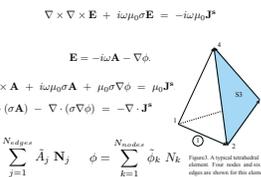


Figure 3. A typical unstructured element. Four nodes and six edges are shown in the diagram.

Verifications and Examples

1 - Controlled-source Electromagnetics (CSEM):

1 - 1 Horizontal Electric Dipole source (HED) for a half-space model

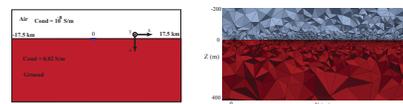


Figure 4. The 3D geometry of the half-space model. The 1 meter electric dipole source is located on the air-ground interface.

The quality mesh is generated using the TetGen Library (Wang, 2007) which uses a Delaunay point distribution of the entire computational domain. The dimension of the whole mesh is 31 km x 31 km x 35 km. The respective number of nodes and edges are 10977 and 19673 which leads to 27268 of elements for the coarse and scalar potentials together. The modeling system is solved using the iterative solver of COMSOL (Computational Modelling Resource) from SPINUS (2009) with a preconditioning of the coefficient matrix. Here after about 1000 iterations a residual norm of 10^{-10} was reached. We also chose a 100-parameter 1000×3 for the ILU-T preconditioner. The dimension of the Kronecker subspace was chosen as 100 for this case.

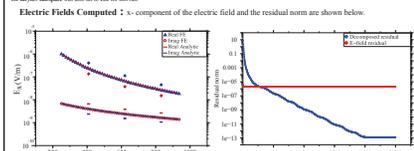


Figure 5. The x-component of the electric field for a frequency of 1 Hz. The blue straight and red circles are the real and imaginary parts of the electric field calculated using the FEM method. The red and blue solid lines are the real and imaginary parts of the electric field calculated using the analytic formula given by Ward and Hobbins (1988).

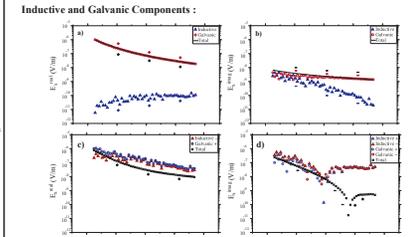


Figure 6. The inductive and galvanic parts of the electric field for frequencies of 1 Hz and 100 Hz for a conductive half-space model of 0.02 S/m. Positive and negative values are shown in red and blue, respectively. (a) Real part of the inductive and galvanic components of the electric field. (b) Imaginary part of the inductive and galvanic components of the electric field. (c) Real part of the inductive and galvanic components of the electric field. (d) Imaginary part of the inductive and galvanic components of the electric field. The location of the source-receiver pair is located at x = 25 km here.

1 - 2 Long Grounded Wire source for a prism in a half-space model

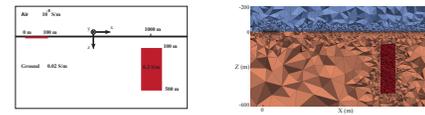


Figure 9. The 3D geometry of the prism model. The 100 meters grounded wire source is located on the air-ground interface. The dimension of the conductive prism is 120 x 200 x 100 in the x, y, and z directions.

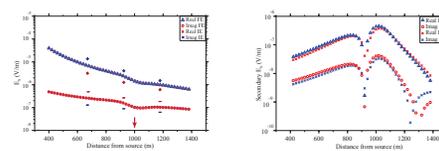


Figure 10. An enlarged x-z cross-section of the central part of the unstructured grid. The conductive ground and the receiver are shown in red and blue respectively. The Earth surface is red and the mesh is refined about the dipole source at the origin (x = 0 and z = 0) and along 1000 meters of observation distance.

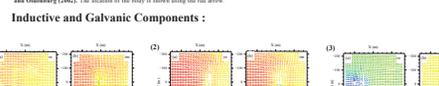


Figure 11. A comparison of the electric field for a frequency of 3 Hz. The blue straight and red circles are the real and imaginary parts of the electric field calculated using the FEM method. The red and blue curves are also the real and imaginary parts of the field calculated using Integral-Equation of Farquharson and Mackenzie (1995). The location of the body is shown using the red arrow.

Inductive and Galvanic Components :

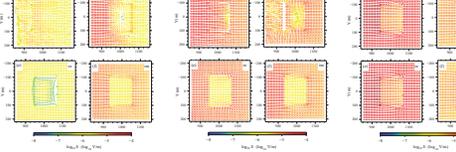


Figure 12. A comparison for the real and imaginary parts of the secondary electric field computed using the finite-element method (blue straight and red circles) and the integral-equation method (red and blue curves).

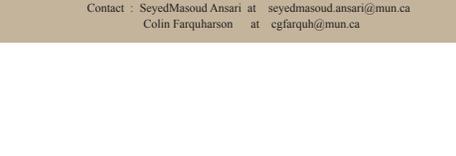


Figure 13. Arrows of the inductive part, the galvanic part and the total electric field for the prism in the half-space model. Fields are shown for three frequencies: (1) 1 kHz, (2) 10 kHz, and (3) 100 kHz. For each frequency all six graphs show the horizontal component of the field at a depth of 10 m in the ground. Direction and strength of the fields are characterized using arrows and colors. The real (a), and the imaginary (b) components of the inductive part, the real (c), and the imaginary (d) component of the galvanic part, and the real (e), and the imaginary (f) components of the total field are shown.

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2 - Transmitter-Receiver pair for a Large Conductivity Contrast

The algorithm is also verified for a model containing a large conductivity contrast. The conductivity ratio of the graphite (shown in red in Figure 14) to the background brine is 8200:1.3. The transmitter-receiver pair, which are two loops of currents, more at a height of 2 cm above the surface of the brine. Here, in order to strongly prove the correctness of our method a comparison with the physical scale modeling (PSM) of Duckworth and Kerbes (1997) and Duckworth et al. (2001) is done. The PSM experiment is performed in a laboratory.

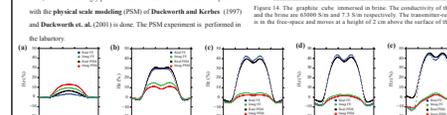


Figure 14. The graphite cube immersed in brine. The conductivity of the graphite and the brine are 82000 S/m and 2.3 S/m respectively. The transmitter-receiver pair is in the free-space and more at a height of 2 cm above the surface of the brine.

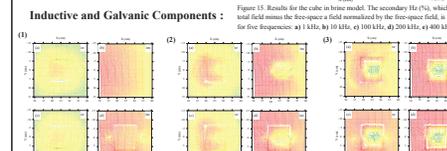


Figure 15. Results for the cube in brine model. The secondary field (V/m) which is the field minus the free-space field normalized by the free-space field, is shown for three frequencies: (a) 1 kHz, (b) 10 kHz, (c) 100 kHz, (d) 200 kHz, (e) 400 kHz.

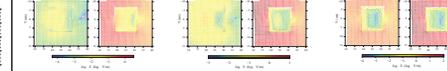


Figure 16. Arrows of the inductive part, the galvanic part and the total electric field for the cube in brine model. Fields are shown for three frequencies: (1) 1 kHz, (2) 10 kHz, and (3) 100 kHz. For each frequency all six graphs show the horizontal component of the field at a depth of 1 cm. Direction and strength of the fields are characterized using arrows and colors. The real (a), and the imaginary (b) components of the inductive part, the real (c), and the imaginary (d) component of the galvanic part, and the real (e), and the imaginary (f) components of the total field are shown. The center of the source-receiver pair is located at x = 25 cm here.

Conclusions

- Unstructured meshes are efficiently used for forward modeling of 3D geophysical electromagnetic problems.
- An $\mathbf{A} - \phi$ decomposition of the electric field leads to a fast solution particularly for lower frequencies.
- The correctness of our algorithm is confirmed for two important geophysical surveys for the commonly used grounded-wire source and for the moving transmitter-receiver loop.
- Inductive and galvanic components are extracted and the interplay between them for different geophysical scenarios are explored.
- The method presented and the algorithm written here are verified for models containing large conductivity contrasts.

References:

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