

Finite volume modelling of electromagnetic data using unstructured staggered grids

Hormoz Jahandari and Colin G. Farquharson



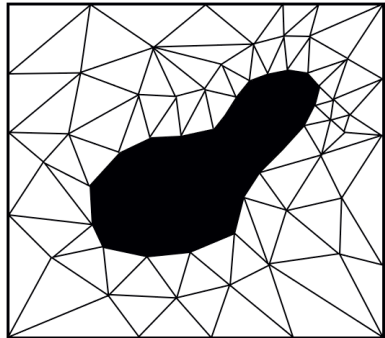
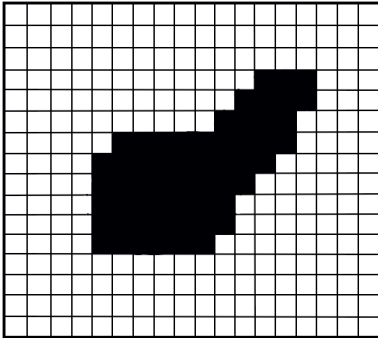
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- 1 Unstructured grids.
- 2 A finite-volume discretization of Maxwell's equations.
- 3 Example for a long grounded wire source.
- 4 Example for magnetic dipole sources.
- 5 Conclusions

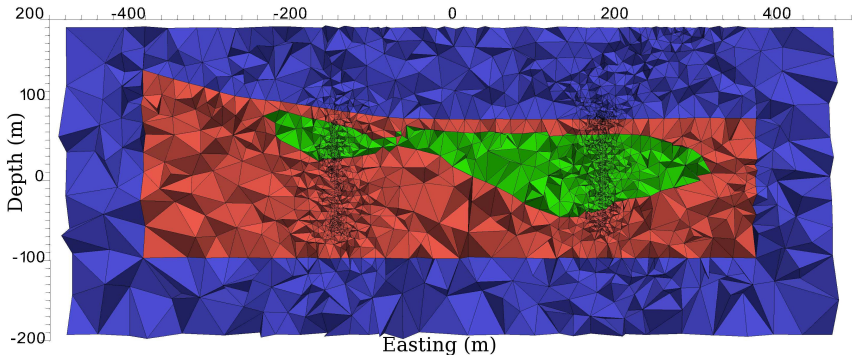
Unstructured grids

- model irregular structures



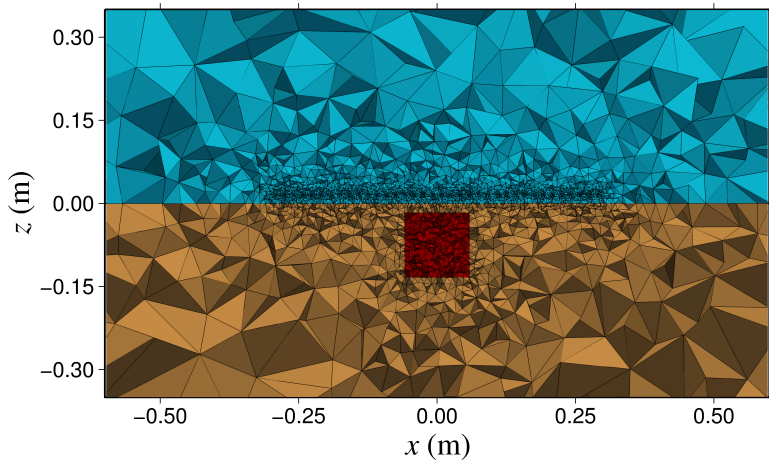
Unstructured grids

- topographical features
- geological interfaces



Unstructured grids

- local refinement (at observation points, sources, interfaces)



Maxwell's equations: describe the behavior of electromagnetic fields.

$$\nabla \times \mathbf{E} = -i\omega\mu_0\mathbf{H} - i\omega\mu_0\mathbf{M}_p$$

$$\nabla \times \mathbf{H} = \sigma\mathbf{E} + \mathbf{J}_p$$

Helmholtz equation for electric field:

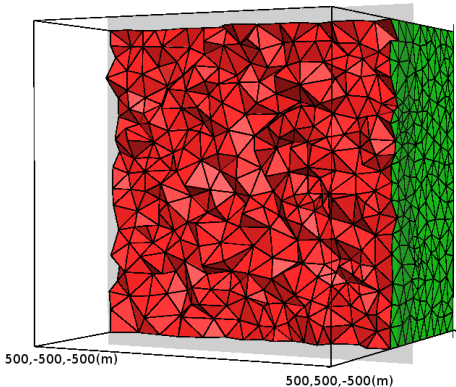
$$\nabla \times \nabla \times \mathbf{E} + i\omega\mu_0\sigma\mathbf{E} = -i\omega\mu_0\mathbf{J}_p - i\omega\mu_0(\nabla \times \mathbf{M}_p)$$

Boundary condition:

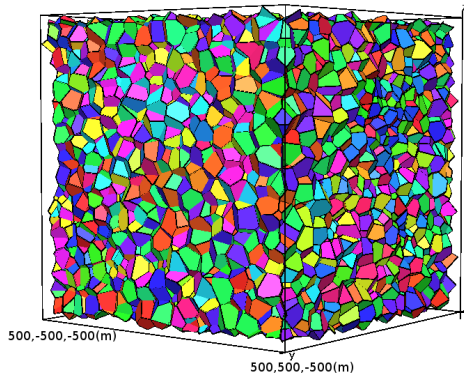
$$\mathbf{E} \cdot \boldsymbol{\tau} = 0 \quad \text{on } \Gamma$$

Staggered tetrahedral-Voronoi grids

tetrahedral grid

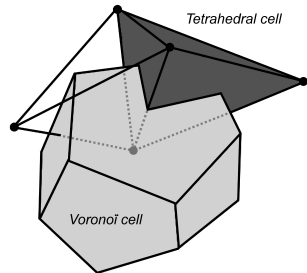
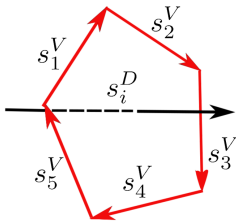
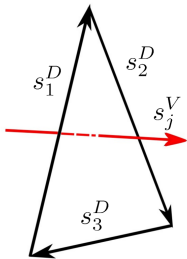


Voronoi grid



Integral form of Maxwell's equations:

$$\oint_{\partial A^D} \mathbf{E} \cdot d\mathbf{l}^D = -i\mu_0\omega \iint_{A^D} \mathbf{H} \cdot d\mathbf{A}^D - i\mu_0\omega \iint_{A^D} \mathbf{M}_p \cdot d\mathbf{A}^D$$
$$\oint_{\partial A^V} \mathbf{H} \cdot d\mathbf{l}^V = \sigma \iint_{A^V} \mathbf{E} \cdot d\mathbf{A}^V + \iint_{A^V} \mathbf{J}_p \cdot d\mathbf{A}^V$$



Discretized Helmholtz equation

Discretized form of Maxwell's equations:

$$\sum_{q=1}^{W_j^D} E_{i(j,q)} I_{i(j,q)}^D = -i\mu_0\omega H_j A_j^D - i\mu_0\omega M_{pj}$$
$$\sum_{k=1}^{W_i^V} H_{j(i,k)} I_{j(i,k)}^V = \sigma E_i A_i^V + J_{pi}$$

Discretized form of Helmholtz equation:

$$\sum_{k=1}^{W_i^V} \left(\left(\sum_{q=1}^{W_j^D} E_{i(j,q)} I_{i(j,q)}^D \right) \frac{I_{j(i,k)}^V}{A_{j(i,k)}^D} \right) + i\omega\mu_0\sigma E_i A_i^V$$
$$= -i\omega\mu_0 \sum_{k=1}^{W_i^V} M_{pj(i,k)} \frac{I_{j(i,k)}^V}{A_{j(i,k)}^D} - i\omega\mu_0 J_{pi}$$

Sparse direct solver: MUMPS (Amestoy et al., 2006)

Interpolation inside tetrahedra: edge vector basis functions

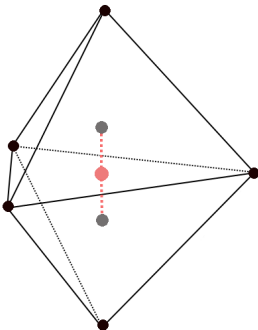
Grid generator: TetGen (Si, 2004)

Inclusion of EM sources

- grounded wire:

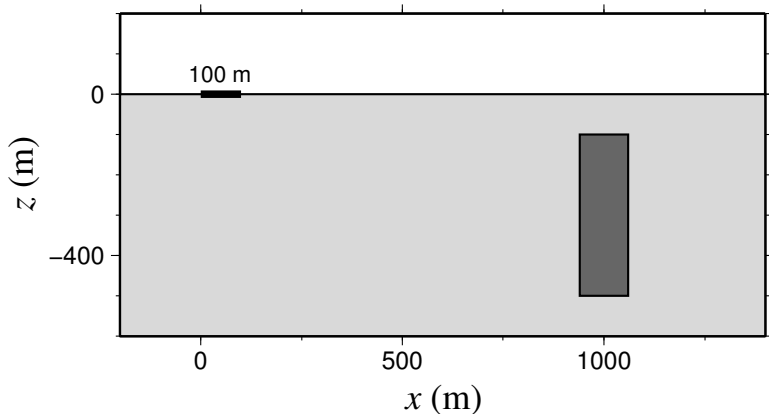


- point vertical magnetic dipole:



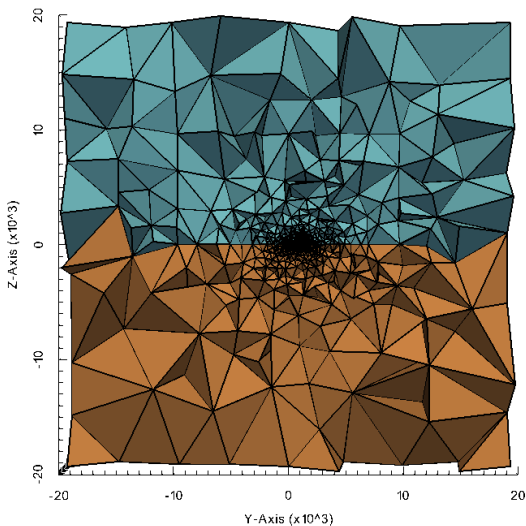
Example 1: long grounded wire

- 100 m wire along the x axis operating at 3 Hz
- dimensions of the prism: $120 \times 200 \times 400$ m
- $\sigma_{ground} = 0.02$ S/m ; $\sigma_{prism} = 0.2$ S/m
- observation points along the x axis



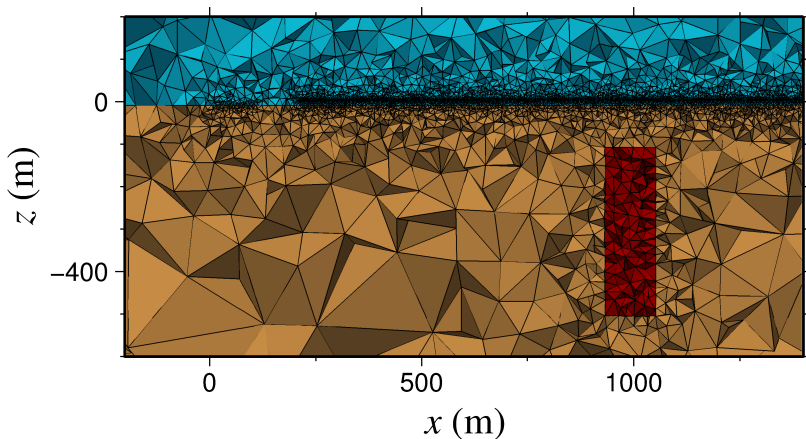
Example 1: long grounded wire

- dimensions of the domain: $40 \times 40 \times 40$ km
- number of tetrahedra: 162,689 ; number of unknowns: 189,105



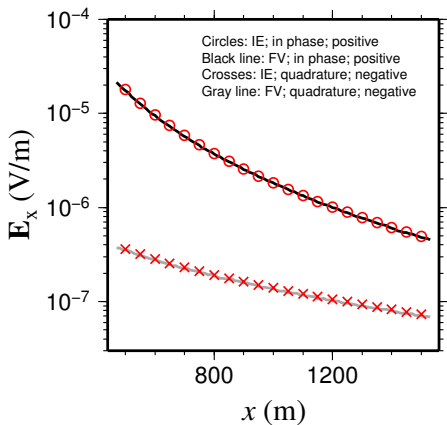
Example 1: long grounded wire

- grid refined at the source, observation points and the prism
- computation time: 40 s ; memory: 4 Gbytes (on Apple Mac Pro; 2.26 GHz Quad-Core Intel Xeon processor)

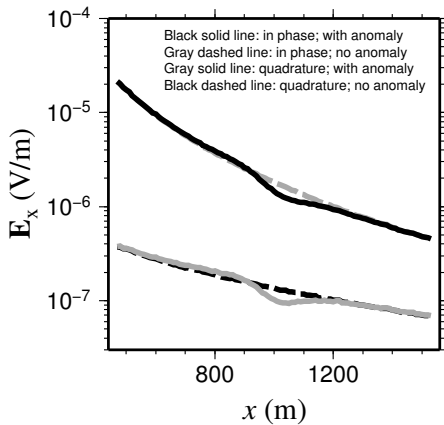


Example 1: long grounded wire

- without prism (homogeneous halfspace)
- total field
- FV vs IE (Farquharson and Oldenburg, 2002)

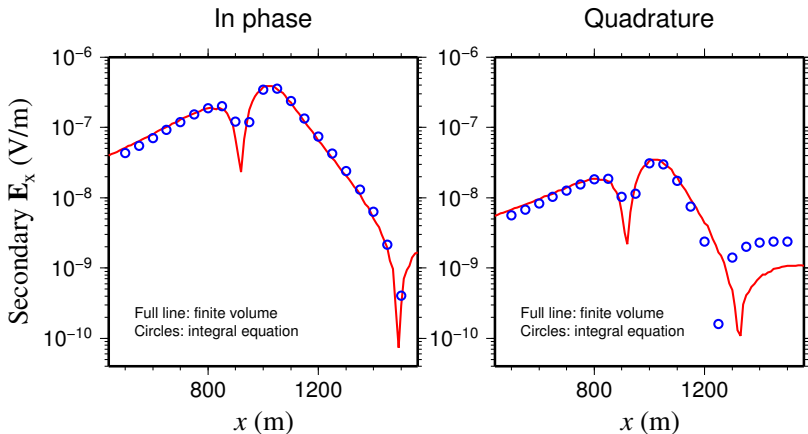


- with and without prism
- total field
- FV only



Example 1: long grounded wire

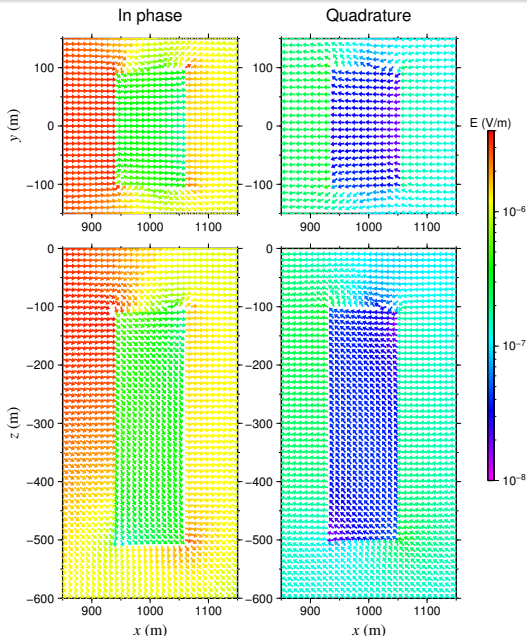
- scattered field
- FV vs IE



Example 1: long grounded wire

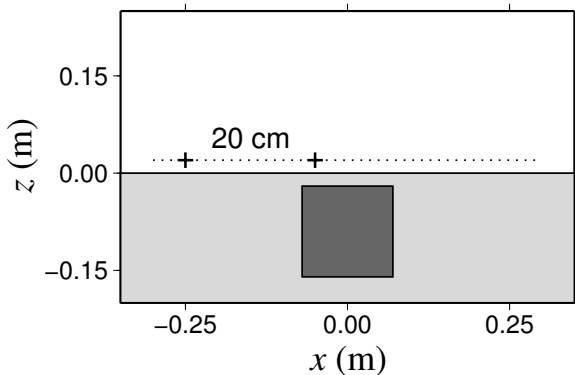
- horizontal section ($z = -150 \text{ m}$)
- total electric field (in phase and quadrature)

- vertical section ($y = 0 \text{ m}$)
- total electric field (in phase and quadrature)



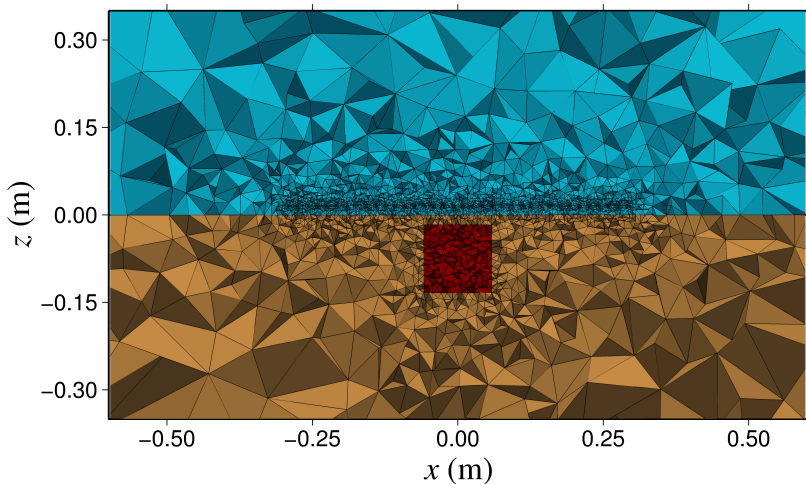
Example 2: magnetic dipole transmitter-receiver pairs

- graphite cube in brine (physical scale modelling measurements)
- transmitter-receiver pairs along the x axis at $z = 2 \text{ cm}$
- dimensions of the cubic graphite: $14 \times 14 \times 14 \text{ cm}$
- $\sigma_{brine} = 7.3 \text{ S/m}$; $\sigma_{prism} = 63,000 \text{ S/m}$
- frequencies: 1, 10, 100, 200, 400 kHz



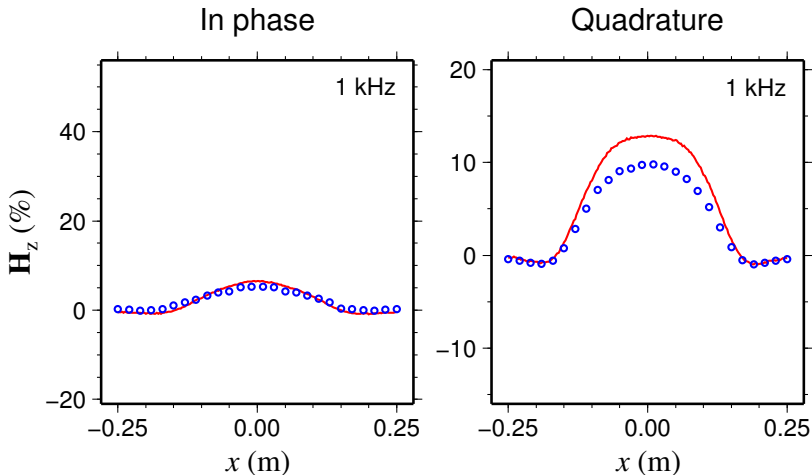
Example 2: magnetic dipole transmitter-receiver pairs

- grid refined at the sources, observation points and the prism



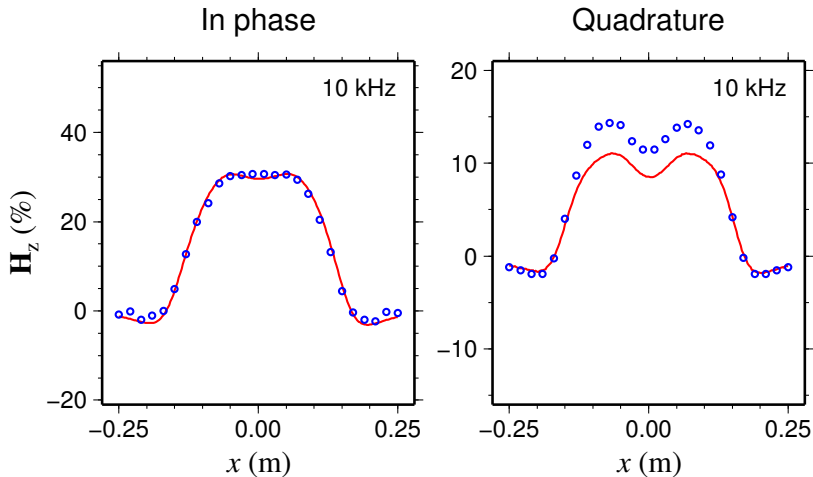
Example 2: magnetic dipole transmitter-receiver pairs

- scattered H-field: (total–free-space)/free-space
- FV vs PSM (Farquharson et al., 2006)



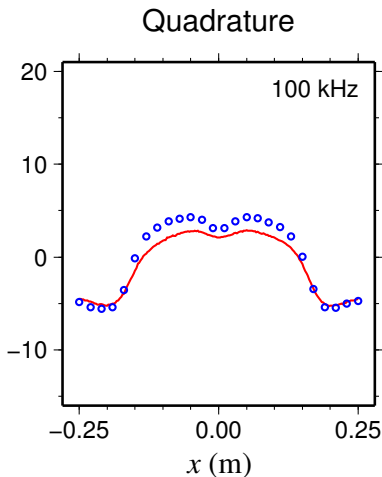
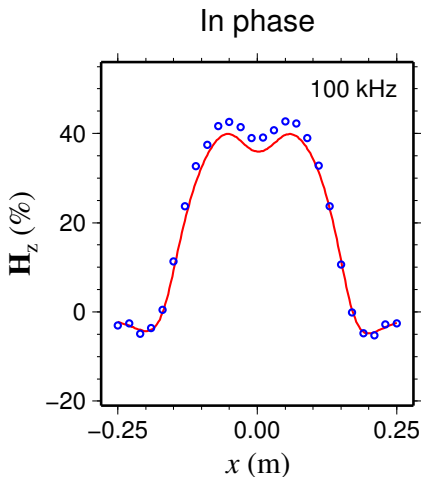
Example 2: magnetic dipole transmitter-receiver pairs

- scattered H-field: (total–free-space)/free-space
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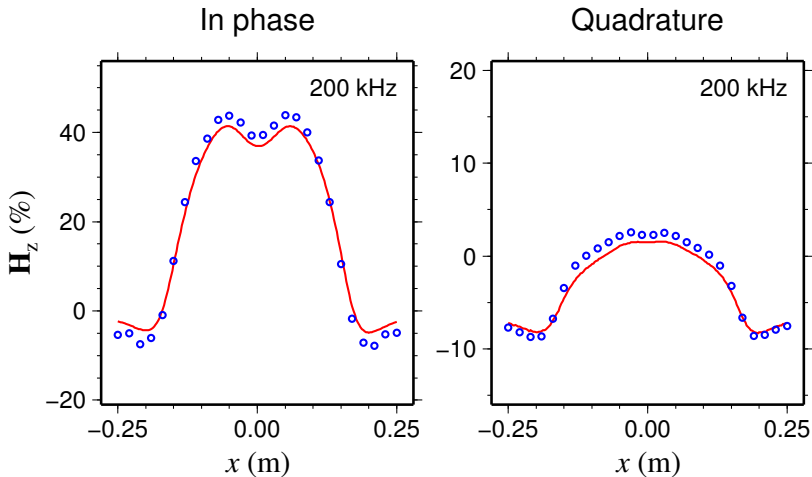
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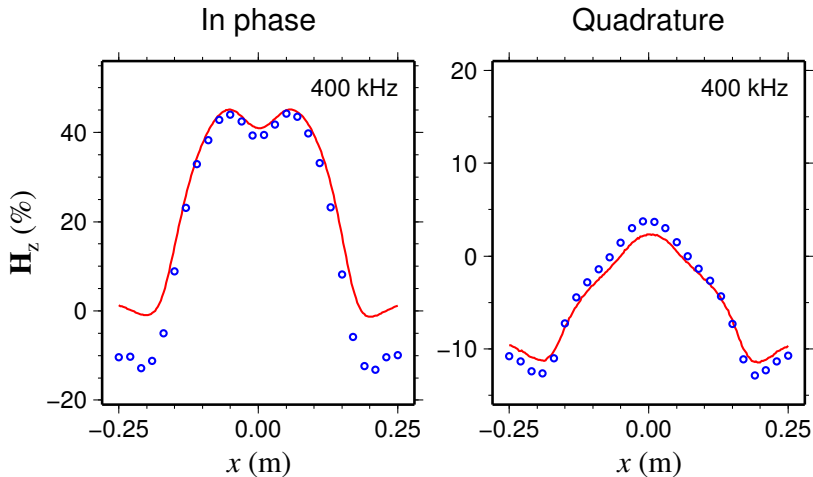
Example 2: magnetic dipole transmitter-receiver pairs

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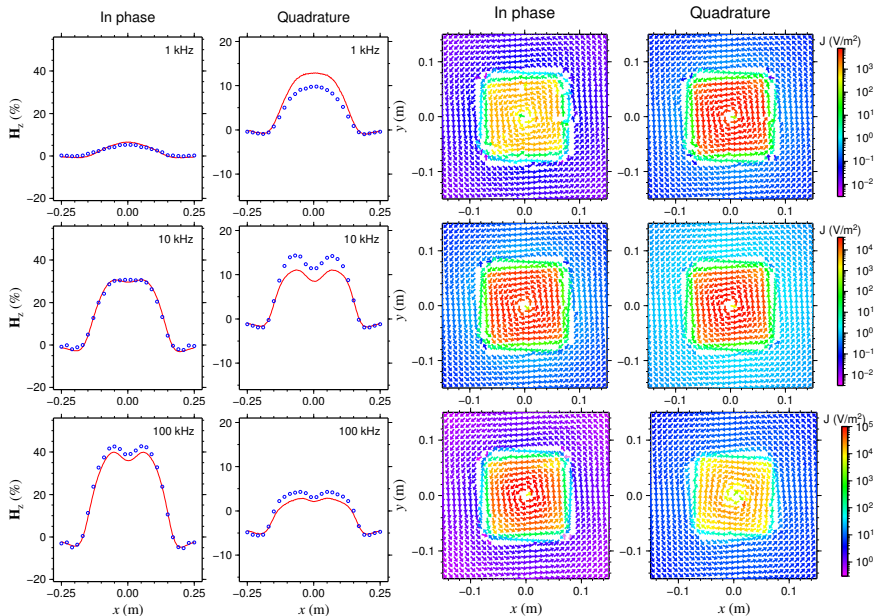


Example 2: magnetic dipole transmitter-receiver pairs

- scattered H-field: (total – free-space)/free-space
- FV vs PSM

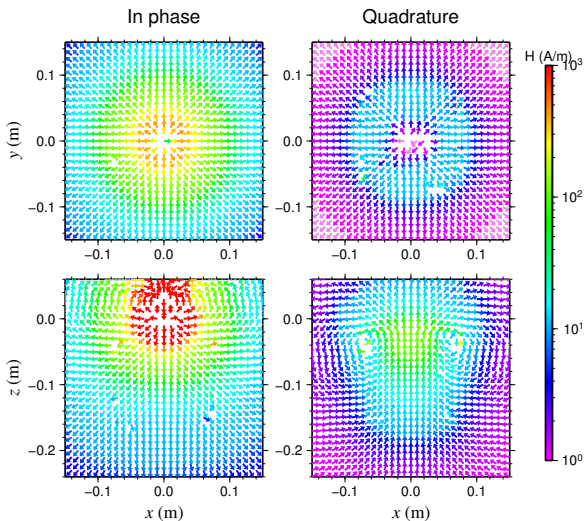


Example 2: magnetic dipole transmitter-receiver pairs



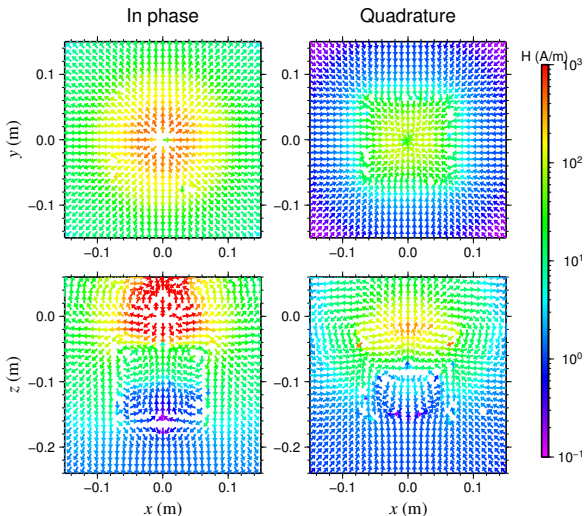
Example 2: magnetic dipole transmitter-receiver pairs

- total H-field (1 kHz)
- horizontal and vertical sections



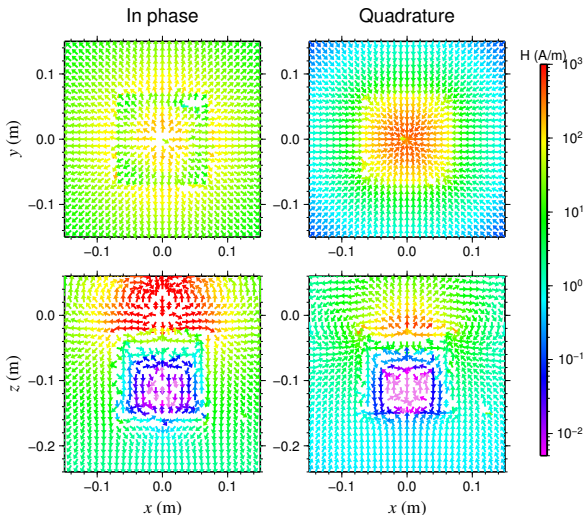
Example 2: magnetic dipole transmitter-receiver pairs

- total H-field (10 kHz)
- horizontal and vertical sections



Example 2: magnetic dipole transmitter-receiver pairs

- total H-field (100 kHz)
- horizontal and vertical sections



- A finite-volume technique is used for modelling the electromagnetic data. This technique uses the staggered tetrahedral-Voronoi grid.
- The Helmholtz equation is discretized and solved using a sparse direct solver (MUMPS).
- The scheme has been tested for two models: one with a long grounded wire source; another one for magnetic source-receiver pairs with large conductivity contrasts.
- For the both examples, the results from the FV scheme are in good agreement with those from the IE method (example 1) and the physical scale modelling measurements (example 2).

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Atlantic Canada
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NSERC
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VALE

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